Flows of SU(2)-structures and beyond

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(joint work with Henrique Sá Earp)

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Want to address this problem for an *H*-structure with $H \subset O(n)$.

My main interest is when H belongs to Berger's list:

Dim	Holonomy group	Parallel tensor	Geometry
2 <i>n</i>	U(<i>n</i>)	ω, J	Kähler
2 <i>n</i>	$\mathrm{SU}(n)$	ω, Ψ	Calabi-Yau
4 <i>n</i>	$\operatorname{Sp}(n)$	$\omega_1, \omega_2, \omega_3$	HyperKähler
4 <i>n</i>	$\operatorname{Sp}(n)\operatorname{Sp}(1)$	Ω	Quaternion-Kähler
7	G_2	φ	G_2
8	Spin(7)	Φ	Spin(7)

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Holonomy reduction is equivalent to $\nabla \eta = 0$, where $\operatorname{Stab}(\eta) \cong H \subset SO(n)$. Interested in finding solutions by means of geometric flow i.e.

$$\frac{\partial}{\partial t}\eta = \cdots$$

One can also consider Sasaki, 3-Sasaki, PSU(3),... structures

General framework for flows of H-structures

Given a 2-tensor $A \in \mathfrak{gl}(\mathfrak{n}, \mathbb{R}) \cong T^* \otimes T^*$ we can act by endomorphism \diamond on any tensor α i.e. $(A \diamond \alpha)_{ij\ldots k} = A_{ia}\alpha_{aj\ldots k} + \cdots + A_{ka}\alpha_{ij\ldots a}$.

[Fadel-Loubeau-Moreno-Sá Earp 22], [Karigiannis 08]

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Note that if $A \in \mathfrak{h}$ then $A \diamond \eta = 0$, so we need

$$A \in \operatorname{End}(\mathbb{R}^n)_{\mathfrak{h}} := \operatorname{End}(\mathbb{R}^n)/\mathfrak{h} = S^2 \oplus \mathfrak{h}^{\perp}.$$

We now know which spaces A has to belong to, but how do we choose it?

Invariant quantities: rep theory

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Bryant 03

Given $H \subset O(n)$, the space of kth order invariants $V_k(\mathfrak{h})$ is implicitly given by

$$V_k(\mathfrak{h}) \oplus (\Lambda^1 \otimes S^{k+1}) = (S^2 \oplus \mathfrak{h}^{\perp}) \otimes S^k,$$

where $S^k = S^k(\Lambda^1)$ and $\mathfrak{h}^{\perp} = \Lambda^2/\mathfrak{h}$.

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Proof.

 $V_1(\mathfrak{h}) = \Lambda^1 \otimes \mathfrak{h}^{\perp}$ is the space of intrinsic torsion. On the other hand $V_2(\mathfrak{h})$ is spanned by ∇T and Rm but these overlap in $\mathfrak{h}^{\perp} \otimes \Lambda^2$ due to a Bianchi type identity: $\nabla T \diamond \eta = \nabla^2 \eta + \text{l.o.t}$. Can iterate this argument for other $V_k(\mathfrak{h})$ using differential Bianchi identity.

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Let us illustrate this explicitly in a few cases!

Cartan, Weyl

 $V_1(\mathfrak{so}(n)) = 0$ $V_2(\mathfrak{so}(n)) = \mathbb{1}\mathbb{R} \oplus \mathbb{1}S_0^2(\mathbb{R}^n) \oplus W$ $\operatorname{End}(\mathbb{R}^n)_{\mathfrak{so}(n)} = \mathbb{1}\mathbb{R} \oplus \mathbb{1}S_0^2(\mathbb{R}^7)$

There are 3 second order invariants: Scalar, Ricci and Weyl curvature.

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1 \times 1 + 1 \times 1 = 2
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There is a 2-parameter of O(n) flows!

$$\partial_t(g) = \lambda_1 \operatorname{Scal}(g)g + \lambda_2 \operatorname{Ric}(g).$$

Bryant 03, Dwivedi-Karigiannis-Gianniotis 23

$$\begin{split} V_1(\mathfrak{g}_2) &= \mathbb{R} \oplus \mathbb{R}^7 \oplus S_0^2(\mathbb{R}^7) \oplus \mathfrak{g}_2 \\ V_2(\mathfrak{g}_2) &= \mathbb{1}\mathbb{R} \oplus 2\mathbb{R}^7 \oplus \mathbb{3}S_0^2(\mathbb{R}^7) \oplus V_{0,1} \oplus 2V_{1,1} \oplus V_{0,2} \oplus V_{3,0} \\ & \text{End}(\mathbb{R}^7)_{\mathfrak{g}_2} = \mathbb{1}\mathbb{R} \oplus \mathbb{1}\mathbb{R}^7 \oplus \mathbb{1}S_0^2(\mathbb{R}^7) \end{split}$$

 $1 \times 1 + 2 \times 1 + 3 \times 1 = 6$

Bryant 03, Dwivedi-Karigiannis-Gianniotis 23

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There is a 6-parameter of G_2 flows to second order!

 $\partial_t(\varphi) = (\lambda_1 \operatorname{Scal}(g)g + \lambda_2 \operatorname{Ric}(g) + \lambda_3 W_{27} + \lambda_4 \mathcal{L}_{\tau_1}g + \lambda_5 \operatorname{div}(T) + \lambda_6 \operatorname{div}(T^t)) \diamond \varphi.$

Recall that $T \in \Lambda^1 \otimes \mathfrak{g}_2^{\perp} \subset \Lambda^1 \otimes \Lambda^2$ is defined by $T \diamond \varphi = \nabla \varphi$.

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Recall that $T \in \Lambda^1 \otimes \mathfrak{g}_2^{\perp} \subset \Lambda^1 \otimes \Lambda^2$ is defined by $T \diamond \varphi = \nabla \varphi$.

Rmk: For $\partial_t(\varphi) = \Delta_{\varphi}\varphi$ AND $d\varphi_t = 0$, $\lambda_2 = -1$ and $\lambda_i = 0$ otherwise.

The Spin(7) case

F20, Dwivedi 24

 $V_1(\mathfrak{spin}(7)) = \mathbb{R}^8 \oplus \mathbb{R}^{48}$ $V_2(\mathfrak{spin}(7)) = \mathbb{1}\mathbb{R} \oplus \mathbb{1}\mathbb{R}^7 \oplus 2S_0^2(\mathbb{R}^8) \oplus V_{0,1,0} \oplus V_{1,1,0} \oplus V_{2,0,0} \oplus V_{0,2,0} \oplus V_{1,0,2}$ $\operatorname{End}(\mathbb{R}^8)_{\mathfrak{spin}(7)} = \mathbb{1}\mathbb{R} \oplus \mathbb{1}\mathbb{R}^7 \oplus \mathbb{1}S_0^2(\mathbb{R}^8)$

 $1 \times 1 + 2 \times 1 + 1 \times 1 = 4$

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There is a 4-parameter of Spin(7) flows to second order!

 $\partial_t(\Phi) = (\lambda_1 \mathrm{Scal}(g)g + \lambda_2 \mathrm{Ric}(g) + \lambda_3 \mathcal{L}_{T_1}g + \lambda_4 \mathrm{div}(T)) \diamond \Phi.$

The SU(3) case

Bedulli-Vezzoni 07

 $V_{1}(\mathfrak{su}(3)) = 2\mathbb{R} \oplus 2\mathbb{R}^{6} \oplus 2\mathfrak{su}3 \oplus \mathbb{R}^{12}$ $V_{2}(\mathfrak{su}(3)) = 3\mathbb{R} \oplus 4\mathbb{R}^{6} \oplus 5\mathfrak{su}3 \oplus 4\mathbb{R}^{12} \oplus 3V_{2,1} \oplus V_{2,2} \oplus V_{3,0}$ $\operatorname{End}(\mathbb{R}^{6})_{\mathfrak{su}(3)} = 2\mathbb{R} \oplus 1\mathbb{R}^{6} \oplus 1\mathfrak{su}3 \oplus 1\mathbb{R}^{12}$

 $3 \times 2 + 4 \times 1 + 5 \times 1 + 4 \times 1 = 19$

There is a 19-parameter of SU(3) flows to second order!

In fact out of the 3 + 4 + 5 + 4 = 16 invariants, 7 are computable only from curvature: 3 of which come from Ricci and 4 from Weyl!

The SU(2) case

F.-Sá Earp 24

$$V_1(\mathfrak{su}(2)) = 3\mathbb{R}^4$$
$$V_2(\mathfrak{su}(2)) = 9\mathbb{R} \oplus 12\mathbb{R}^3 \oplus \mathbb{R}^5$$
$$\operatorname{End}(\mathbb{R}^4)_{\mathfrak{su}(2)} = 4\mathbb{R} \oplus 3\mathbb{R}^3$$

 $9 \times 4 + 12 \times 3 = 72$

There is a 72-parameter of SU(2) flows to second order! Let's be more explicit in this case.

The SU(2) case

F.-Sá Earp 24

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Rmk: $V_2(\mathfrak{sp}(2)) = 13\mathbb{R} \oplus 27\mathfrak{sp}(2) \oplus 24\mathbb{R}^5 \oplus \cdots$ so it gets harder in higher dimension....

First order invariants

The intrinsic torsion $T \in \Lambda^1 \otimes \mathfrak{su}(2)^{\perp} \subset \Lambda^1 \otimes \Lambda^2$ is defined by

 $T(\cdot)\diamond\boldsymbol{\omega}=\nabla_{\cdot}\boldsymbol{\omega},$

where $\boldsymbol{\omega} \coloneqq (\omega_1, \omega_2, \omega_3)$, but $\mathfrak{su}(2)^{\perp} \cong \Lambda^2_+$, so we have

$$T = \mathbf{a}_1 \otimes \omega_1 + \mathbf{a}_2 \otimes \omega_2 + \mathbf{a}_3 \otimes \omega_3 \in \mathbf{3}\mathbb{R}^4 \cong V_1(\mathfrak{su}(2))$$

Explicitly,

$$a_{1} = \frac{1}{2} (-* d\omega_{1} - J_{2} * d\omega_{2} + J_{3} * d\omega_{3})$$

$$a_{2} = \frac{1}{2} (-* d\omega_{2} - J_{3} * d\omega_{3} + J_{1} * d\omega_{1})$$

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First order invariants

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where $\boldsymbol{\omega} \coloneqq (\omega_1, \omega_2, \omega_3)$, but $\mathfrak{su}(2)^{\perp} \cong \Lambda^2_+$, so we have

$$T = \mathbf{a_1} \otimes \omega_1 + \mathbf{a_2} \otimes \omega_2 + \mathbf{a_3} \otimes \omega_3 \in \mathbf{3R}^4 \cong V_1(\mathfrak{su}(2))$$

Explicitly,

$$a_{1} = \frac{1}{2}(-*d\omega_{1} - J_{2} * d\omega_{2} + J_{3} * d\omega_{3})$$

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Next we want: Second order invariants $V_2(\mathfrak{su}(2))!$

Second order invariants

Recall that

$$V_2(\mathfrak{su}(2)) = {}^{9}\mathbb{R} \oplus {}^{12}\mathbb{R}^3 \oplus \mathbb{R}^5.$$

F-Sá Earp

All second-order invariants of an ${\rm SU}(2)\mbox{-structure}$ can be expressed in terms of the following linearly independent terms:

- The 9 functions $\delta a_1, \delta a_2, \delta a_3$ and $g(da_i, \omega_j)$, where $1 \le i \le j \le 3$.
- **2** The 12 anti-selfdual forms $\pi^2_-(da_i)$, where i = 1, 2, 3, and $S^2_{0,j}(\nabla a_i)$, where i, j = 1, 2, 3.
- **3** The anti-selfdual Weyl curvature $W^- \in S_0^2(\Lambda_-^2)$.

In particular, this implies that Ric and W^+ are expressible in terms of (1) and (2) + lower order terms in $S^2(V_1(\mathfrak{su}(2)))$.

$$W^{+} = -\frac{1}{24} \Big(2g(da_{1}, \omega_{1}) - g(da_{2}, \omega_{2}) - g(da_{3}, \omega_{3}) \\ + g(J_{3}a_{1}, a_{2}) - g(J_{2}a_{1}, a_{3}) - 2g(J_{1}a_{2}, a_{3}) \Big) (2\omega_{1}^{2} - \omega_{2}^{2} - \omega_{3}^{2}) \\ - \frac{1}{8} \Big(g(da_{2}, \omega_{2}) - g(da_{3}, \omega_{3}) + g(J_{3}a_{1}, a_{2}) + g(J_{2}a_{1}, a_{3}) \Big) (\omega_{2}^{2} - \omega_{3}^{2}) \\ - \frac{1}{4} \Big(g(da_{2}, \omega_{1}) + g(J_{1}a_{1}, a_{3}) \big) (\omega_{1} \odot \omega_{2}) \\ - \frac{1}{4} \Big(g(da_{3}, \omega_{2}) + g(J_{2}a_{2}, a_{1}) \big) (\omega_{2} \odot \omega_{3}) \\ - \frac{1}{4} \Big(g(da_{1}, \omega_{3}) + g(J_{3}a_{3}, a_{2}) \big) (\omega_{3} \odot \omega_{1}) \in S_{0}^{2}(\Lambda_{+}^{2}).$$

Explicit formulae: Ricci

$$\begin{split} \mathrm{Scal}(g) &= -2 \Big(\sum_{i=1}^{3} g(da_{i}, \omega_{i}) + g(J_{1}(a_{1}), J_{2}(a_{2})) + g(J_{1}(a_{1}), J_{3}(a_{3})) \\ &+ g(J_{2}(a_{2}), J_{3}(a_{3})) \Big). \end{split}$$

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where recall $P_{i} : \Lambda^{2}_{-} \xrightarrow{\simeq} S^{2}_{0,i}$ and
$$\Phi_{1} := da_{1} + \frac{1}{4}(J_{2}(a_{2}) \wedge J_{2}(a_{3}) - J_{1}(a_{1}) \wedge J_{2}(a_{3}) + J_{2}(a_{2}) \wedge a_{1}) \\ &- \frac{1}{4}(J_{1}(a_{1}) \wedge J_{3}(a_{2}) + J_{3}(a_{3}) \wedge a_{1}) - \frac{3}{4}(J_{3}(a_{3}) \wedge J_{3}(a_{2})), \end{aligned}$$

and similar expressions hold for Φ_2, Φ_3 .

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where recall $P_{i} : \Lambda_{-}^{2} \xrightarrow{\simeq} S_{0,i}^{2}$ and
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and similar expressions hold for Φ_2, Φ_3 .

Let's go back to $\mathop{\rm SU}(2)\text{-flows}$

Any $\mathrm{SU}(2)$ -flow can be expressed as

$$\partial_t(\omega) = A \diamond \omega$$

where $A = f_0g + B_{0,1} + B_{0,2} + B_{0,3} + f_1\omega_1 + f_2\omega_2 + f_3\omega_3$

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where $A = f_0g + B_{0,1} + B_{0,2} + B_{0,3} + f_1\omega_1 + f_2\omega_2 + f_3\omega_3$

$$\partial_t \omega_1 = + f_0 \omega_1 + f_3 \omega_2 - f_2 \omega_3 + \tilde{B}_{0,1}$$
$$\partial_t \omega_2 = -f_3 \omega_1 + f_0 \omega_2 + f_1 \omega_3 + \tilde{B}_{0,2}$$
$$\partial_t \omega_3 = + f_2 \omega_1 - f_1 \omega_2 + f_0 \omega_3 + \tilde{B}_{0,3}$$

Recall that there is a 72-parameter of SU(2)-flows.

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$$\partial_t(\boldsymbol{\omega}) = A \diamond \boldsymbol{\omega}$$

where $A = f_0g + B_{0,1} + B_{0,2} + B_{0,3} + f_1\omega_1 + f_2\omega_2 + f_3\omega_3$

$$\partial_t \omega_1 = +\mathbf{f}_0 \omega_1 + \mathbf{f}_3 \omega_2 - \mathbf{f}_2 \omega_3 + \tilde{B}_{0,1}$$
$$\partial_t \omega_2 = -\mathbf{f}_3 \omega_1 + \mathbf{f}_0 \omega_2 + \mathbf{f}_1 \omega_3 + \tilde{B}_{0,2}$$
$$\partial_t \omega_3 = +\mathbf{f}_2 \omega_1 - \mathbf{f}_1 \omega_2 + \mathbf{f}_0 \omega_3 + \tilde{B}_{0,3}$$

Recall that there is a 72-parameter of SU(2)-flows. But there are some flows which are more natural than other namely those that arise gradient of some functionals.

$$S^{2}(V_{1}(\mathfrak{su}(2))) = \mathbf{15}\mathbb{R} \oplus 21\mathbb{R}^{3}$$

So there are 15 quadratic functionals in T. These are explicitly given by

$$\int_M g(a_i, J_j a_k) \operatorname{vol},$$

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So there are 15 quadratic functionals in T. These are explicitly given by

$$\int_M g(a_i, J_j a_k) \operatorname{vol},$$

and include $\int |T|^2 \text{ vol and } \int |N_J|^2 \text{ vol.}$

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So there are 15 quadratic functionals in T. These are explicitly given by

$$\int_M g(a_i, J_j a_k) \operatorname{vol},$$

and include $\int |T|^2 \text{ vol and } \int |N_J|^2 \text{ vol.}$

By comparison:

$$S^{2}(V_{1}(\mathfrak{g}_{2})) = 4\mathbb{R} \oplus \cdots$$
$$S^{2}(V_{1}(\mathfrak{spin}(7))) = 2\mathbb{R} \oplus \cdots$$
$$S^{2}(V_{1}(\mathfrak{su}(3))) = 11\mathbb{R} \oplus \cdots$$

Rmk: One can also consider functionals coming the \mathbb{R} components of $V_2(\mathfrak{h})$ e.g. $\int_M \operatorname{Scal}(g)$ vol

Consider the 2-parameter family given by

$$\partial_t(\boldsymbol{\omega}) = (-2\operatorname{Ric}(g) - \lambda_1 \mathcal{L}_V g + \lambda_2 \sum_{i=1}^3 \operatorname{div}(a_i)) \diamond \boldsymbol{\omega} =: \mathcal{Q}(\boldsymbol{\omega}),$$

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- If $\lambda_1 = 1$ and $\lambda_2 = 1/2$ then this is negative gradient flow of $\int_M |T|^2$ vol.

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The $\mathop{\rm SU}(2)\text{-flow}$

$$\partial_t(\boldsymbol{\omega}) = (-2\operatorname{Ric}(g) - \lambda_1 \mathcal{L}_{\sum_{i=1}^2 J_i a_i} g + \lambda_2 \sum_{i=1}^3 \operatorname{div}(a_i)) \diamond \boldsymbol{\omega}$$

has short time existence and uniqueness for $9 - 4\sqrt{5} < \lambda_2 + \frac{\lambda_1}{2} < 9 + 4\sqrt{5}$.

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Upshot: Symmetric part is Ricci DeTurck flow (which is strongly parabolic!), so suffices to linearise skewsymmetric part and compute the condition when principal symbol of $L_{\mathcal{P}}$ is invertible: this gives the range.

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S4: Uniqueness: given 2 solutions, we can pullback by the diffeomorphisms generated by λV_i to get 2 Ricci like flows. Use harmonic map flow to get another pair of diffeomorphisms and show that pullback flow is parabolic (hence unique). Then since V_i depend on ω_i , one can argue that they induce the same diffeomorphism.

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- We can add terms involving W⁺ (i.e. g(da_i, ω_j)) into the above flow and still get existence! Are there good choices?
- ² The Lie derivative term can more generally be set to $\sum_{i,j} \lambda_{ij} \mathcal{L}_{J_i a_j} g$
- Questions: What l.o.t to choose? Are there flows preserving torsion classes (or combinations of those) e.g. Laplacian flow? Long time existence, singularities, ...

Thank you for listening!